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Zakład Fizyki Doświadczalnej



Nieskończenie mała abstrakcja

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SYMETRIE

W mechanice kwantowej trzy najważniejsze symetrie związane są z niezmienniczością:

1. na odwrócenie czasu T (*time*)
2. na zmianę ładunku elektrycznego C (*charge*)
3. parzystości P (*parity*)

Ad. 1. Odwrócenie czasu to *puszczenie filmu od tyłu*

Ad. 2. Zmiana ładunku to zastąpienie wszystkich cząstek ich odpowiednikami z *antymaterii* (na przykład protonu antyprotonem, elektronu pozytonem)

Ad. 3. Zmiana parzystości oznacza obserwację układu w hipotetycznym "trójwymiarowym" lustrze

Wszechświat NIE jest niezmienniczy względem przekształceń T, C, P

W elektrodynamice i chemii istnieje symetria parzystości :
obiekty lewoskrętne i prawoskrętne podlegają tym samym prawom.

W biologii i fizyce słabych oddziaływań symetria parzystości jest złamana:
obiekty lewoskrętne i prawoskrętne zachowują się inaczej.

PARZYSTOŚĆ PRZESTRZENNA P

Transformacja parzystości P jest dyskretną transformacją współrzędnych czasoprzestrzeni

$$P: \Phi(x,y,z,t) \rightarrow \Phi(-x,-y,-z,t)$$

Transformacje te tworzą dyskretną grupę $Z_2 = \{P, 1\}$

- 1 opisuje transformację tożsamościową, $P \cdot P = 1$

W mechanice kwantowej transformacji tej towarzyszy istnienie operatora parzystości P . P jest to operatorem unitarnym. Z właściwości grupy wynika, że $P \psi_\lambda = \lambda \psi_\lambda$.

Funkcje własne o określonej parzystości spełniają równanie własne $P \psi_\lambda = \lambda \psi_\lambda$ dla $\lambda^2 = 1$.

Każdemu polu kwantowemu można więc przypisać wielkość fizyczną, którą nazywa się parzystością. Parzystość może być równa -1 lub $+1$.

Stany z parzystością -1 nazywamy stanami *lewoskrętnymi*
a stany z parzystością $+1$ stanami *prawoskrętnymi*.

PARZYSTOŚĆ ŁADUNKOWA C

W obrazie cząstek elementarnych transformacja C zamienia wszystkie cząstki na ich antycząstki.

Antymateria, układ antycząstek.

Antycząstki to cząstki elementarne podobne do występujących w zwykłej materii, ale o przeciwnym znaku ładunku elektrycznego oraz wszystkich addytywnych liczb kwantowych np. izospinu, dziwności, liczby barionowej.

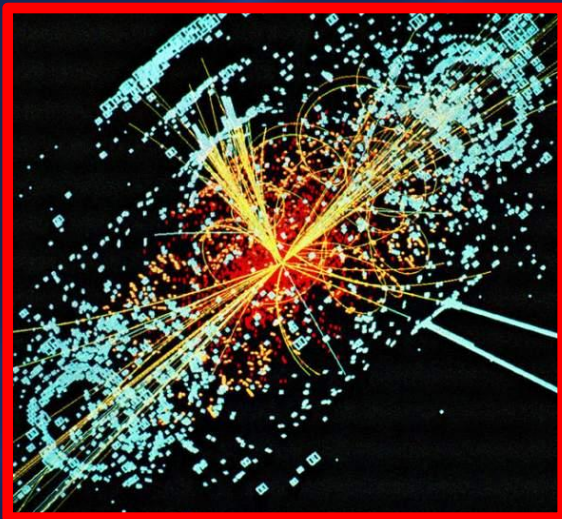
W momencie kontaktu antymaterii ze zwykłą materią, zachodzi anihilacja.

Energia związana z masą spoczynkową anihilujących cząstek ulega przy tym zamianie na energię promieniowania elektromagnetycznego lub energię kinetyczną lżejszych cząstek.

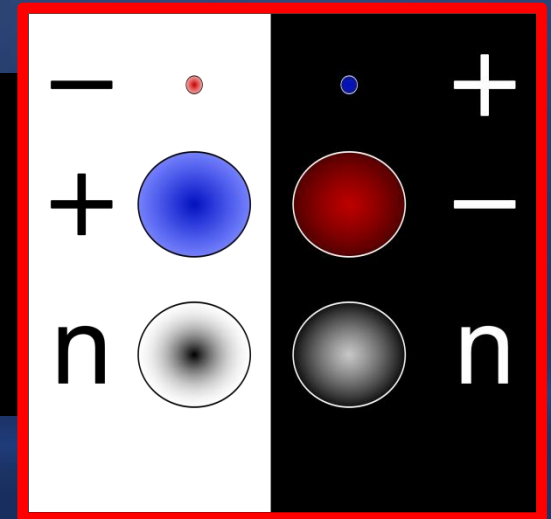
SPONTANICZNE ŁAMANIE SYMETRII

Nagroda Nobla z fizyki za 2008 roku została przyznana **Yoichiro NAMBU** z **Uniwersytetu w Chicago**, **Makoto KOBAYASHI** z **Ośrodka Wysokich Energii w Japonii** i **Toshihide MASKAWA** z **Uniwersytetu w Kyoto** za ich prace na temat łamania symetrii.

Nambu otrzymał nagrodę za wyjaśnianie **mechanizmu spontanicznego łamania symetrii** a Kobayashi i Maskawa za prace teoretyczne wiążące łamanie symetrii z istnieniem w przyrodzie trzech rodzin kwarków.



LHCb
Large Hadron Collider beauty
Detektor cząstek elementarnych
Wielki Zderzacz Hadronów
Genewa, CERN.



Zjawisko łamania CP

ma miejsce przy rozpadach **kwarków dziwnych (s)** na swoich lżejszych kuzynów

(Nie)znany elektron

You know, it would be sufficient to really understand the electron

Albert Einstein

Electric Dipol Moment (EDM)

The electron that can be told is not the true electron

David Harrison

of the electron

During an exam (taken from a book):

Professor: What is an electron?

Student: Ah, God damn it! I have forgotten.

And in fact even in the morning I knew it.

Professor: You should recollect it without fail, Professor said, because you were the unique person who knew what electron was, and you had suddenly forgotten!

Tell me what an electron is and I'll then tell you everything

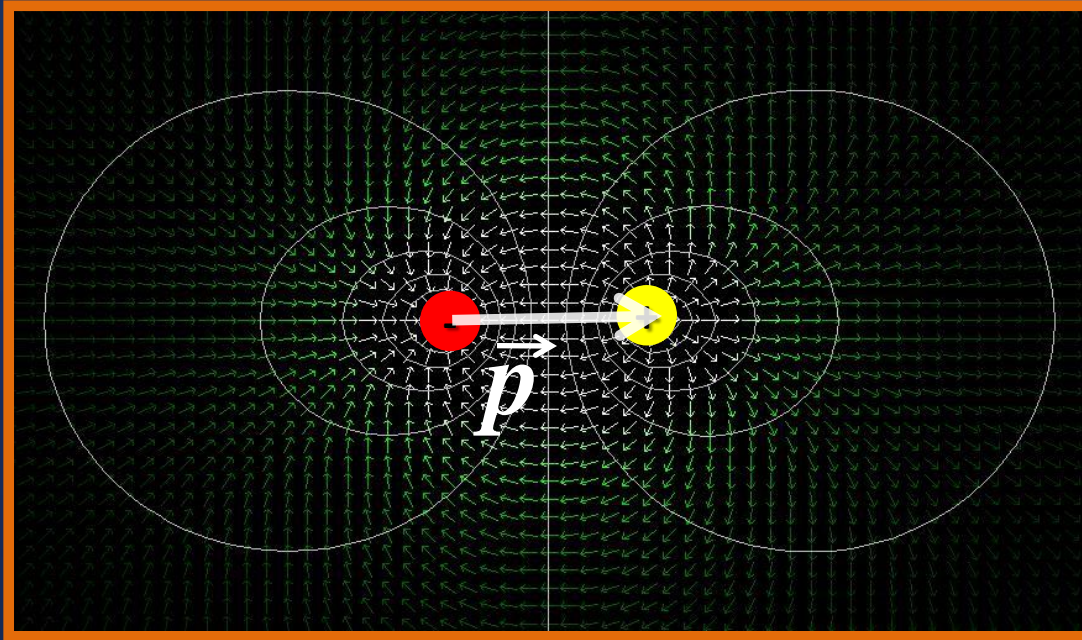
Somebody

Measuring (almost) zero

Chad ORZEL, *Physics World*, Vol. 22, No.12, December 2009

An abstraction ?

Inset: Electric Dipole Moment

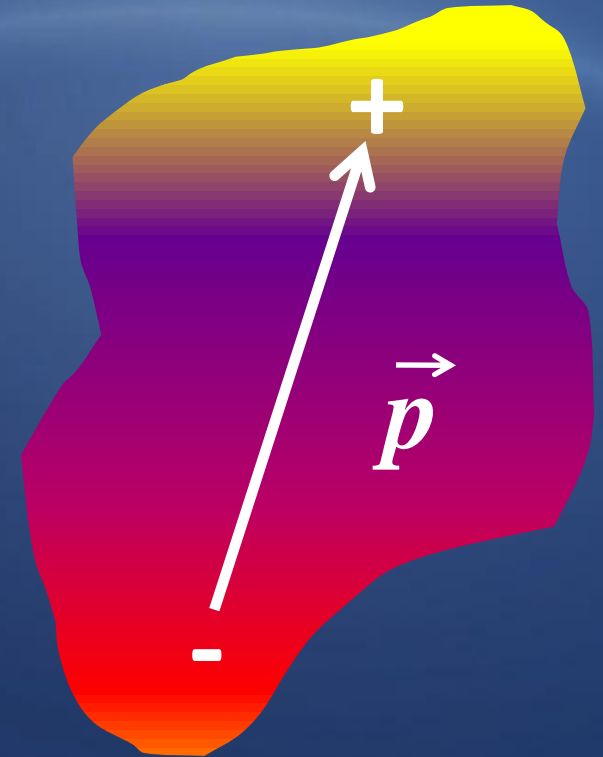


$$\vec{p} = q \cdot \vec{l}$$

p – dipol moment

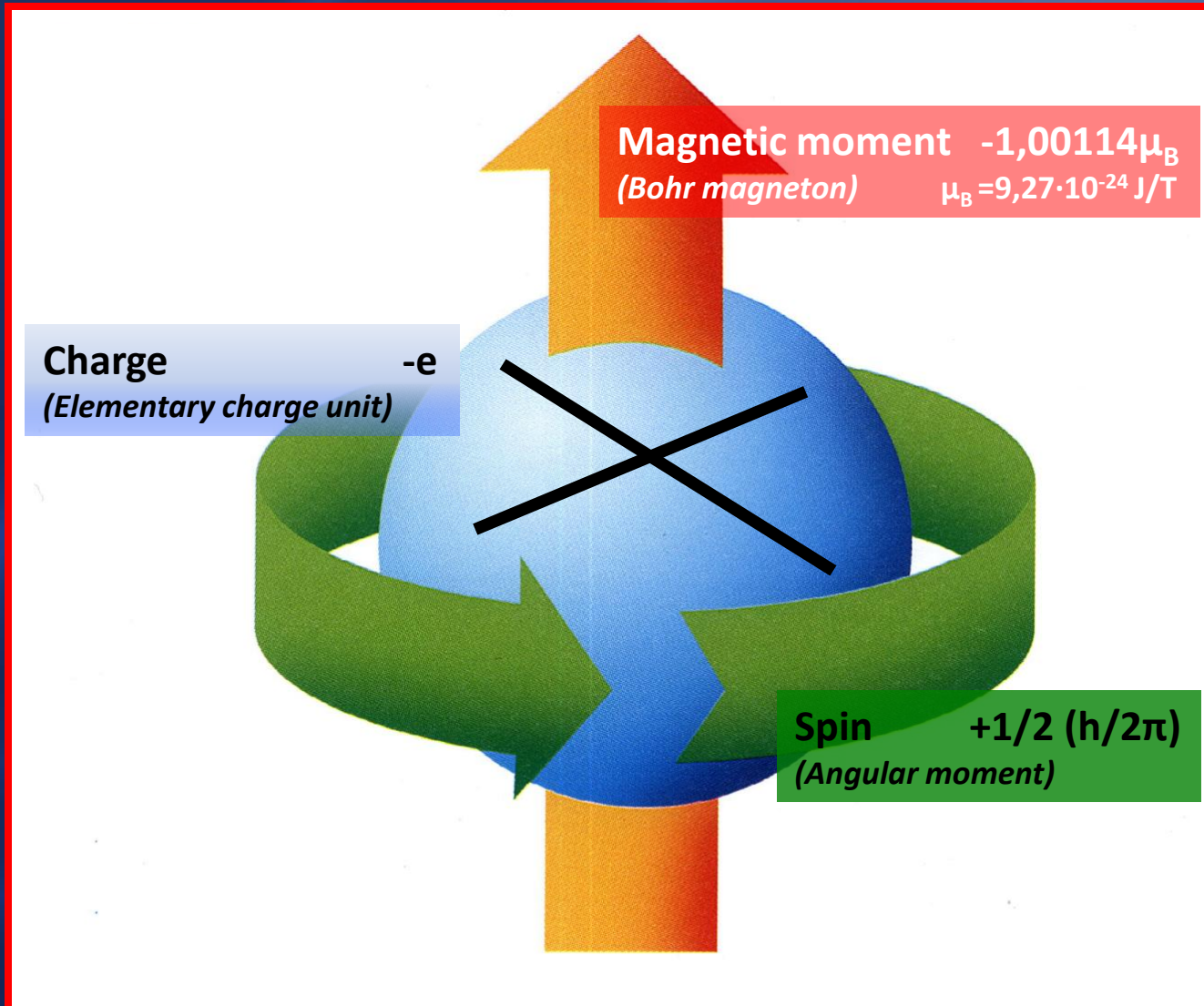
q – value of charge

l – distance between the charges  and 

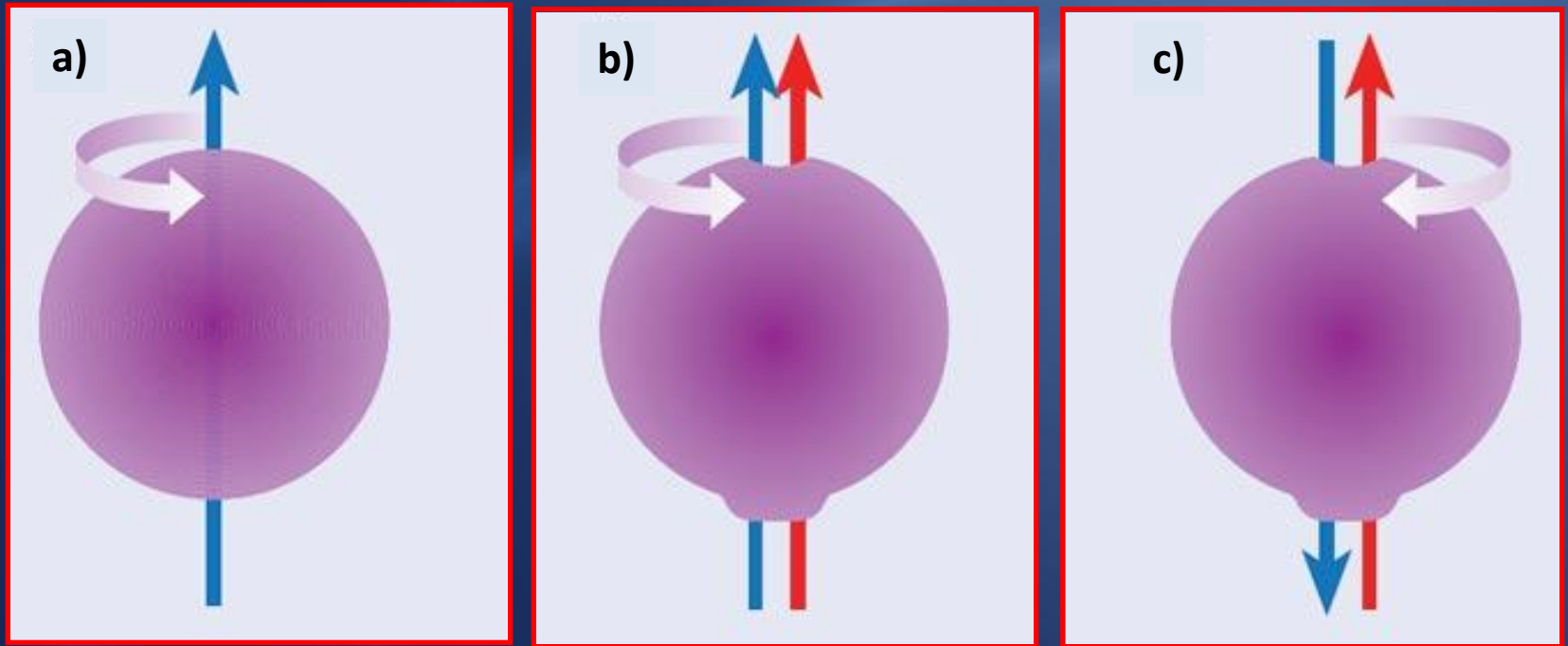


Moment dipol występuje w cząsteczkach w przypadku nierównomiernego rozmieszczenia ładunku dodatniego i ujemnego.

ELECTRON now



Time - reversal violation with an electron EDM



a)

A spinning charge acts like a loop of current to produce a **magnetic dipole moment**

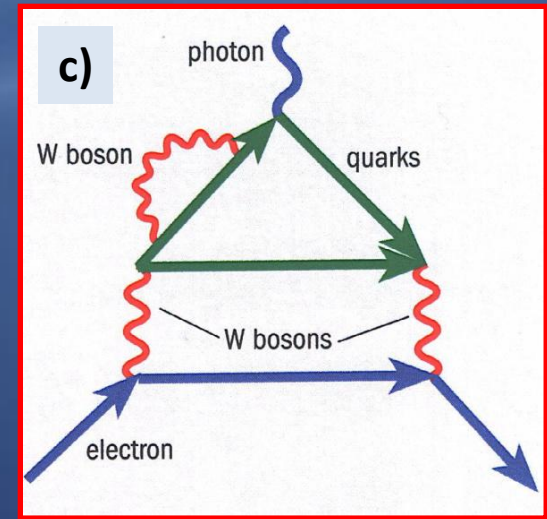
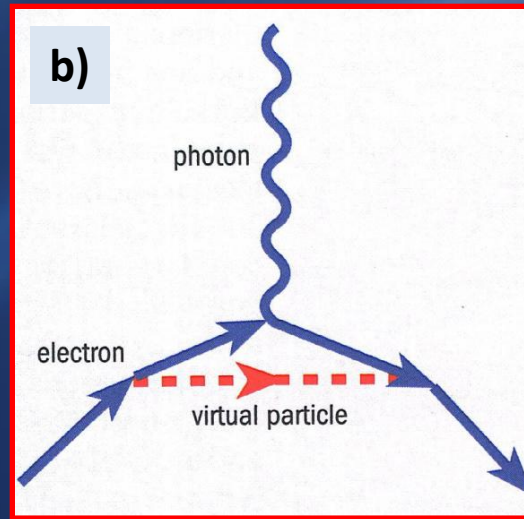
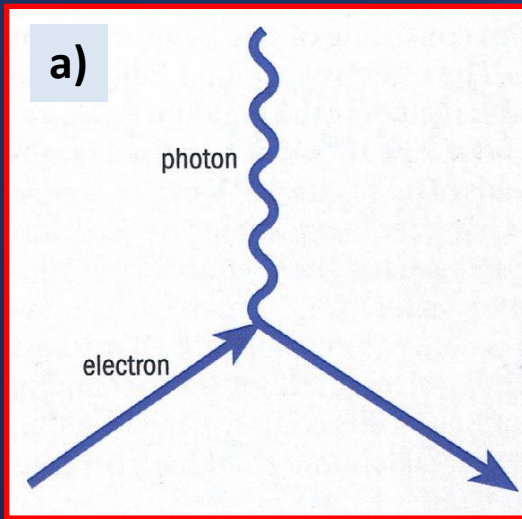
b)

Distorting the electron's charge distribution produces an **electric dipole moment** in the same direction

c)

If time reverses, then the charge spins in the opposite direction. This flips the magnetic dipole but not the electric one, so the two dipoles are no longer aligned – violating time reversal symmetry and preventing the electron from having a permanent electric dipole moment in the basic Standard Model of particle physics.

Time - reversal violation with an electron EDM



a)

The electron-field interaction is described in terms "Feynman diagrams", pictures showing as the electron interacts with its environment. Time flows from left to right, so the illustration shows an electron moving through space, then absorbing or emitting a photon from the electromagnetic field and changing its direction.

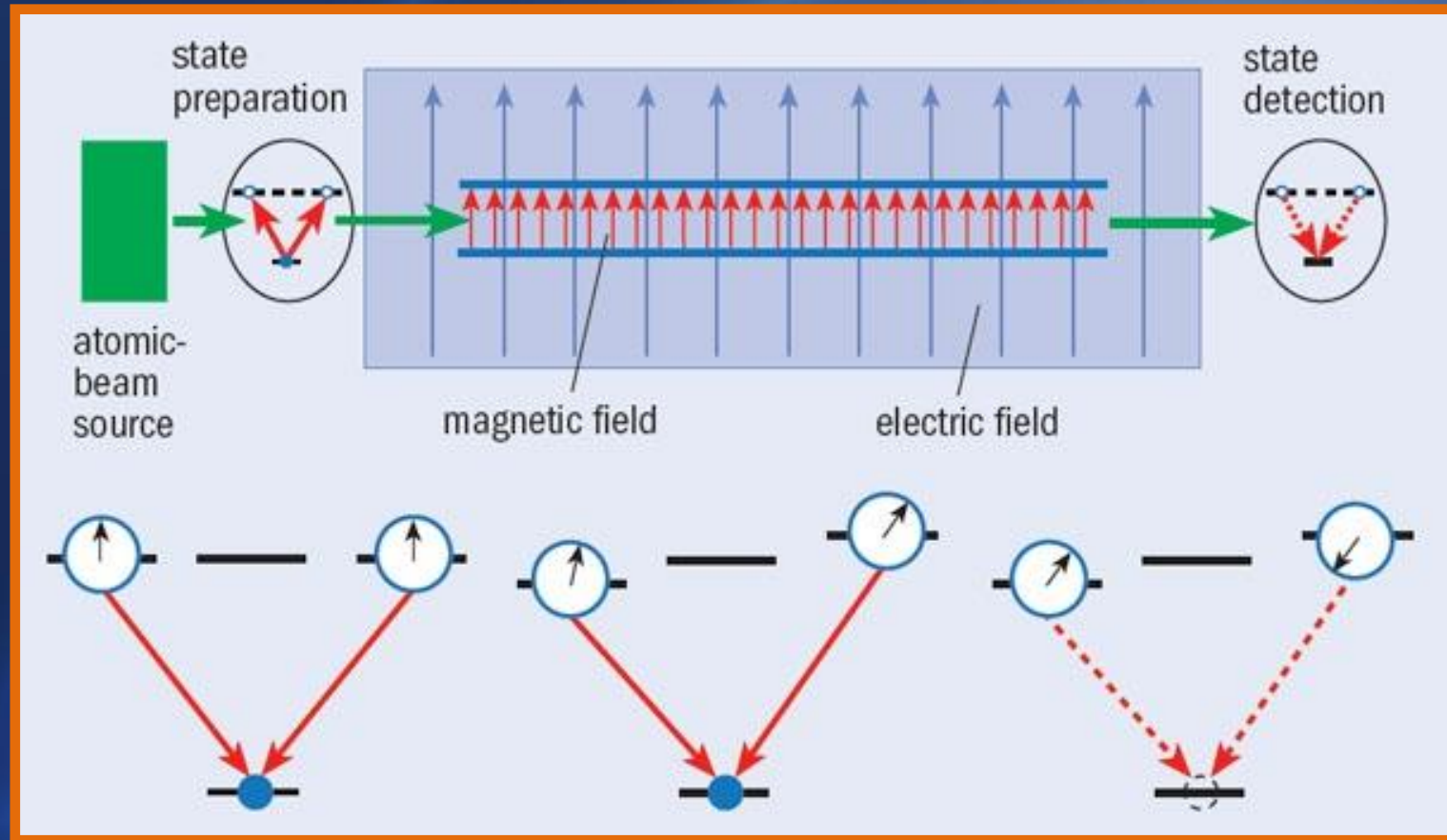
b)

The most important type of the diagram. The electron moves through space as before, but before interacting with the photon, it emits a „virtual“ particle, which it re-absorbs after interacting with the photon. The virtual particle can be just about anything that interacts with an electron, provided it disappears quickly enough that it cannot be detected directly. Such "loop,, diagrams change the way that the electron interacts with light, and an electron electric dipole moment (EDM) is the result. Most theories of physics beyond the Standard Model predict the existence of new particle that can create an EDM through simple diagram like that shown in Figure b).

c)

Within the Standard Model, it is much more difficult to generate an EDM. The simplest diagram that produces an EDM involves three bosons (red lines) and a part of quarks (green lines) in addition to the photon and electron. Such a complicated process is extremely unlikely, and thus produces a very small EDM.

Interferometric detection of an EDM shift

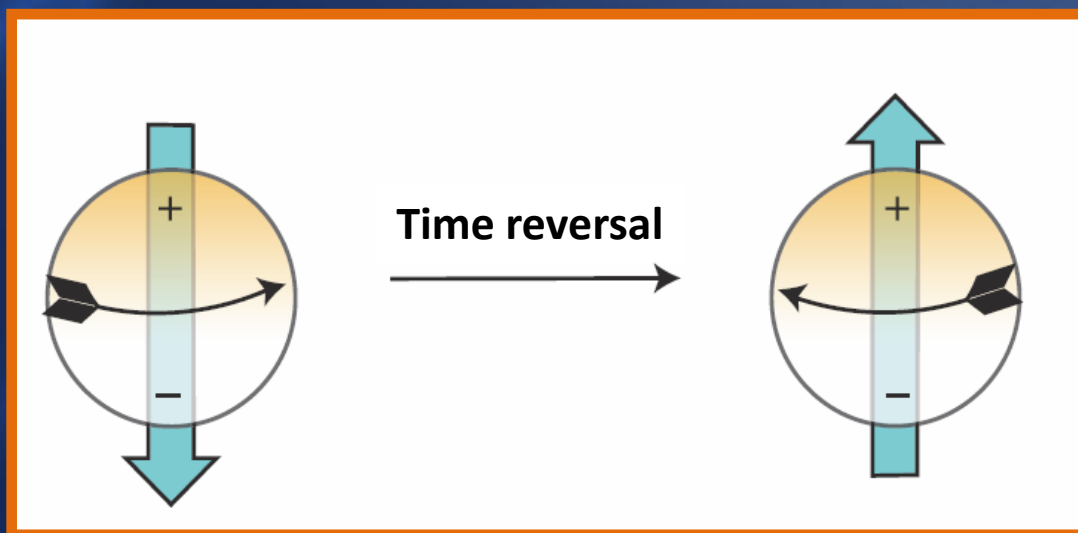


The basic structure of a beam-based electric dipole moment (EDM) experiment.

A multiferroic material to search for the permanent electric dipole moment of the electron

K. Z. Rushchanskii¹, S. Kamba², V. Goian², P. Vaněk², M. Savinov², J. Prokleška³, D. Nuzhnyy², K. Knížek², F. Laufek⁴, S. Eckel⁵, S. K. Lamoreaux⁵, A. O. Sushkov⁵, M. Ležaić¹ and N. A. Spaldin^{6*}

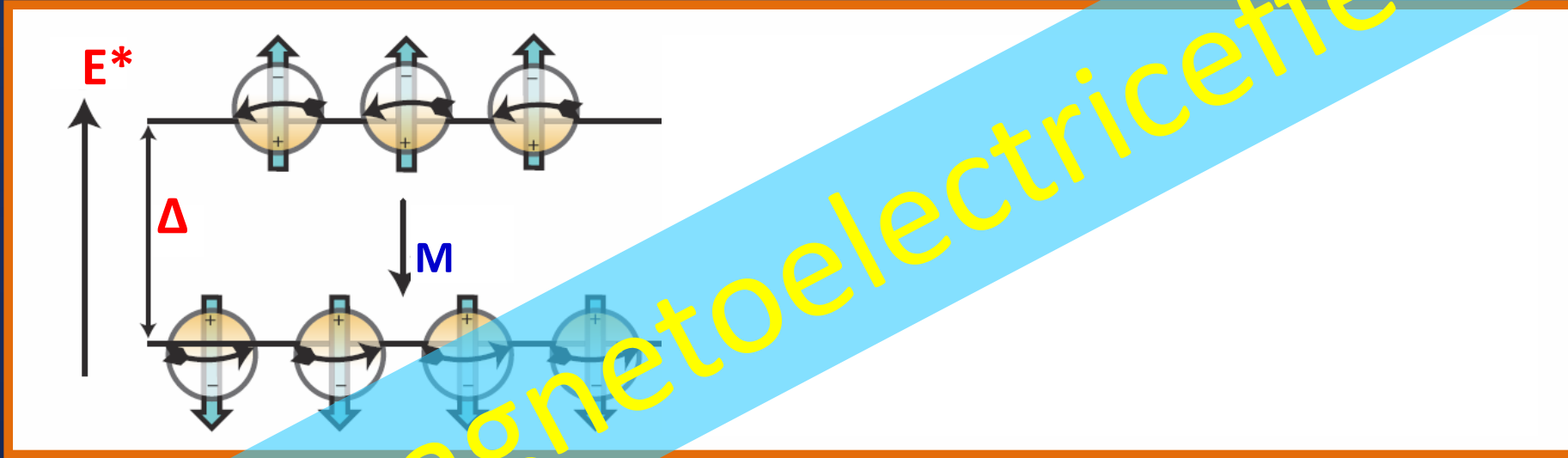
Illustration that an electron with an EDM violates time-reversal symmetry



Both the **EDM** (+ and - symbols; **orange shading**) and **magnetic moment** (blue arrow) of the electron lie along the same axis as the electron spin (black arrow).

The operation of time reversal reverses the magnetic moment but does not affect the EDM. Therefore, an electron with a non-zero EDM violates time-reversal symmetry.

Schematic of the physics underlying the experiment to search for the electron EDM



1. As spin is the only intrinsic vector associated with the electron, a non-vanishing electron EDM is either parallel or antiparallel to its spin and hence its magnetic moment.

2. The energy of electrons with spin up is higher than that for electrons with anti-parallel spin.

$$P_i = \sum \alpha_{ij} H_j + \sum \beta_{ijk} H_j H_{kj} + \dots$$

3. When an electric field - which is parallel and antiparallel to it - is applied, the two electron populations generates a net magnetization.

$$M_i = \sum \alpha_{ij} E_j + \sum \beta_{ijk} E_j E_{kj} + \dots$$

4. When the electric field is reversed there is a magnetization reversal, ΔM , which can be detected using a sensitive magnetometer.

The following material specifications will enable a sensitive EDM search to be mounted:

1. The material should be ferroelectric, with a large electric polarization, and switchable at liquid-He temperature.
2. There should be a high concentration of heavy ions with local magnetic moments that remain paramagnetic at liquid-He temperature; both long-range order and freezing into a glassy state must be avoided.
3. The local environment at each magnetic ion should be strongly modified by the ferroelectric switching.
4. The sample should be macroscopic.

*Superconducting Quantum
Interference Device
Sensitivity $5 \cdot 10^{-18} \text{ T}$*

With these materials properties, and optimal SQUID noise levels, the projected experimental sensitivity is $d_e = 10^{-28} \text{ [e}\cdot\text{cm]}$ after ten days of averaging

ładunek elektronu $e = 1.6 \cdot 10^{-19} \text{ [C]}$

To search for the electron EDM
a specific material is needed



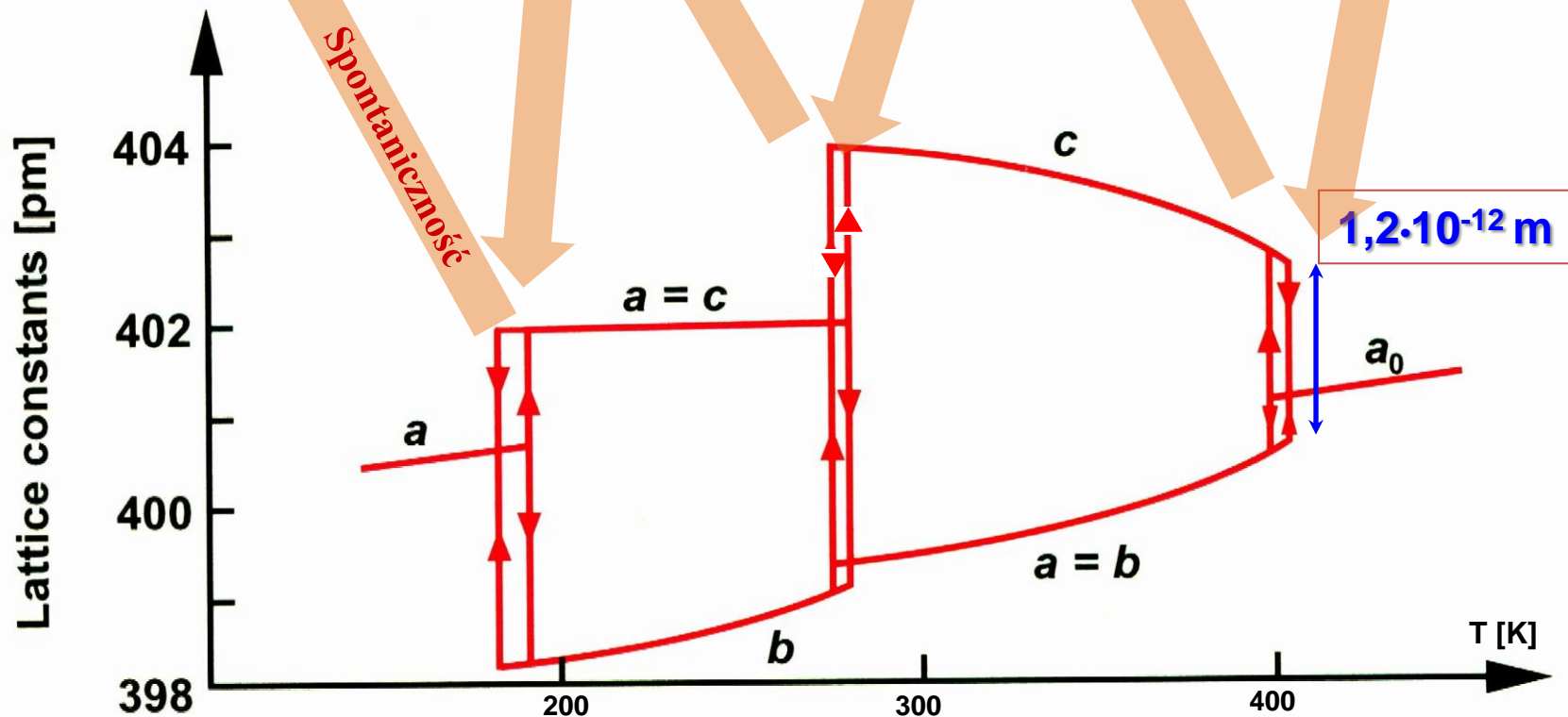
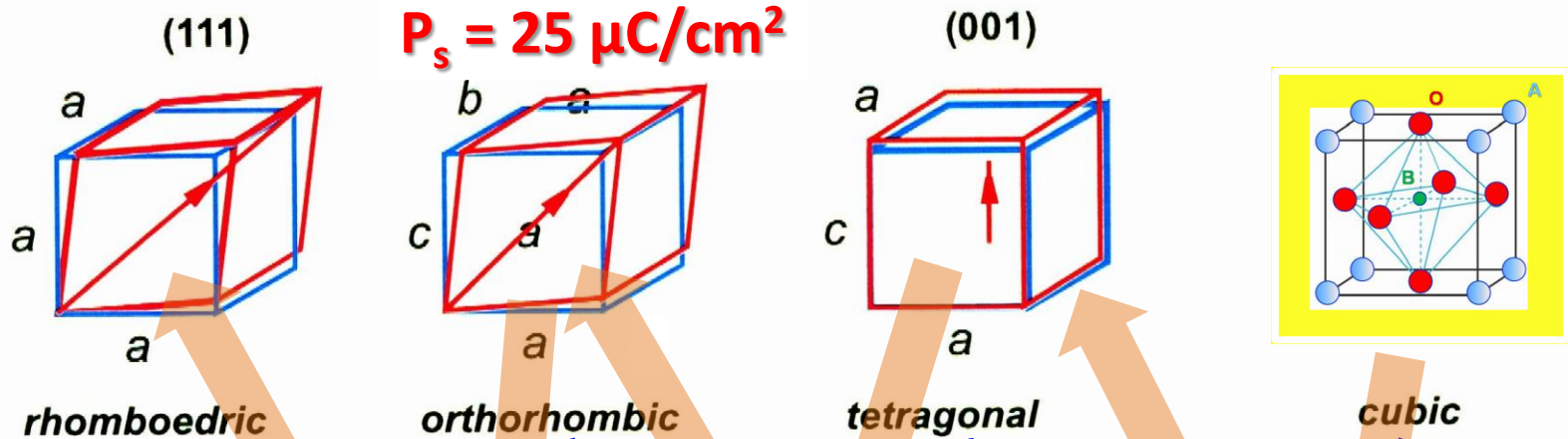
The lattice parameter of EuTiO₃ is 3.905 Å notably smaller than that of BaTiO₃. It is not ferroelectric, but has a large dielectric constant (400) at low temperature indicative of proximity to a ferroelectric phase transition; indeed, it has recently been reported to be a quantum paraelectric.

The Eu⁺² ion has seven unpaired localized 4f electrons, resulting in a large spin magnetization of 7μ_B, and EuTiO₃ is an antiferromagnet with a low Néel temperature of 5.3 K.

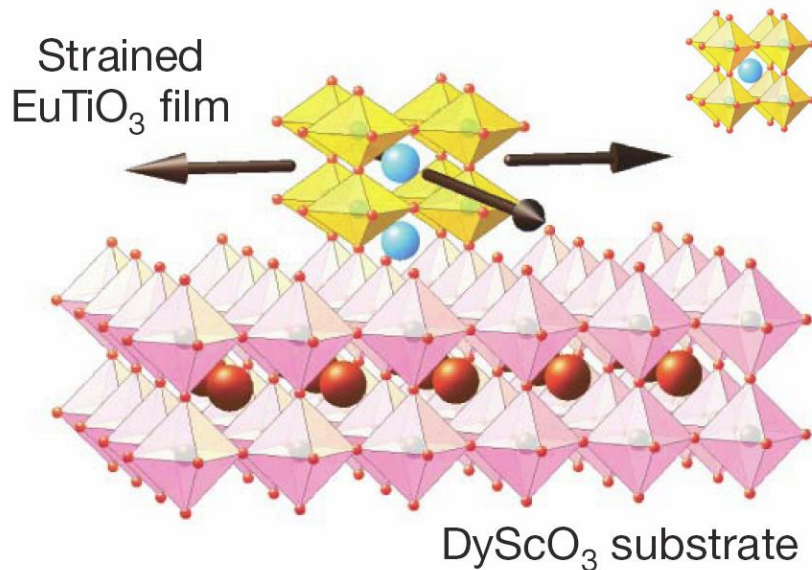
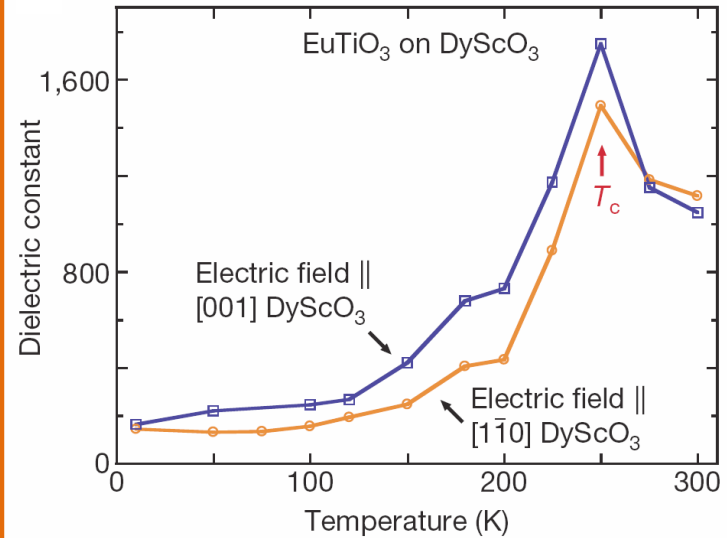
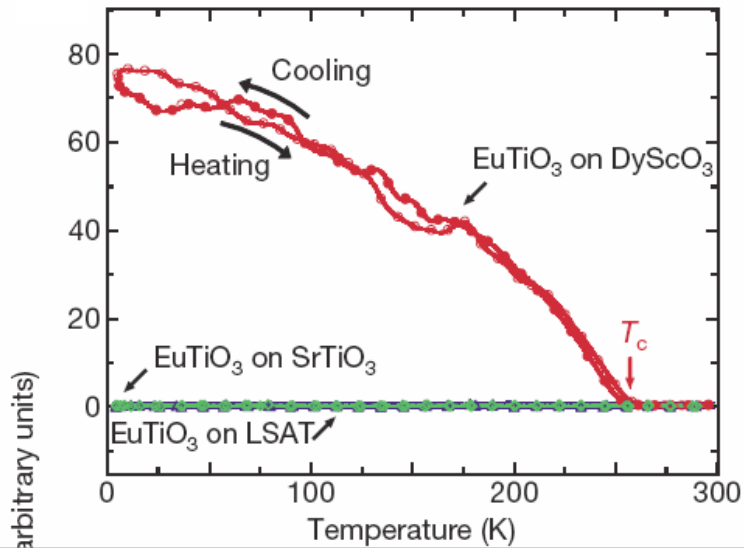
A hypothesis is that by alloying Ba on the A site of EuTiO₃ the magnetic ordering temperature will be suppressed through dilution, and the tendency to ferroelectricity will be increased through the expansion of the lattice constant.

The magnetic ordering temperature is sufficiently low while the ferroelectric polarization and the concentration of magnetic ions remain sufficiently large.

Phase transitions in barium titanate BaTiO₃



Phase transitions in europe titanate EuTiO_3



EuTiO_3 , was predicted to exhibit strong ferromagnetism with spontaneous magnetization, **7 Bohr magnetons per Eu** and strong **ferroelectricity with spontaneous polarization $10 \mu\text{C}/\text{cm}^2$** simultaneously under large biaxial compressive strain. These values are orders of magnitude higher than those of any known ferroelectric ferromagnet and rival the best materials that are solely ferroelectric or ferromagnetic.

Thickness of the EuTiO_3 film is of the order of 20 nm

The best candidate



Eu

Ionic radius: **0.95 Å**

Ba

Ionic radius: **1,35 Å**

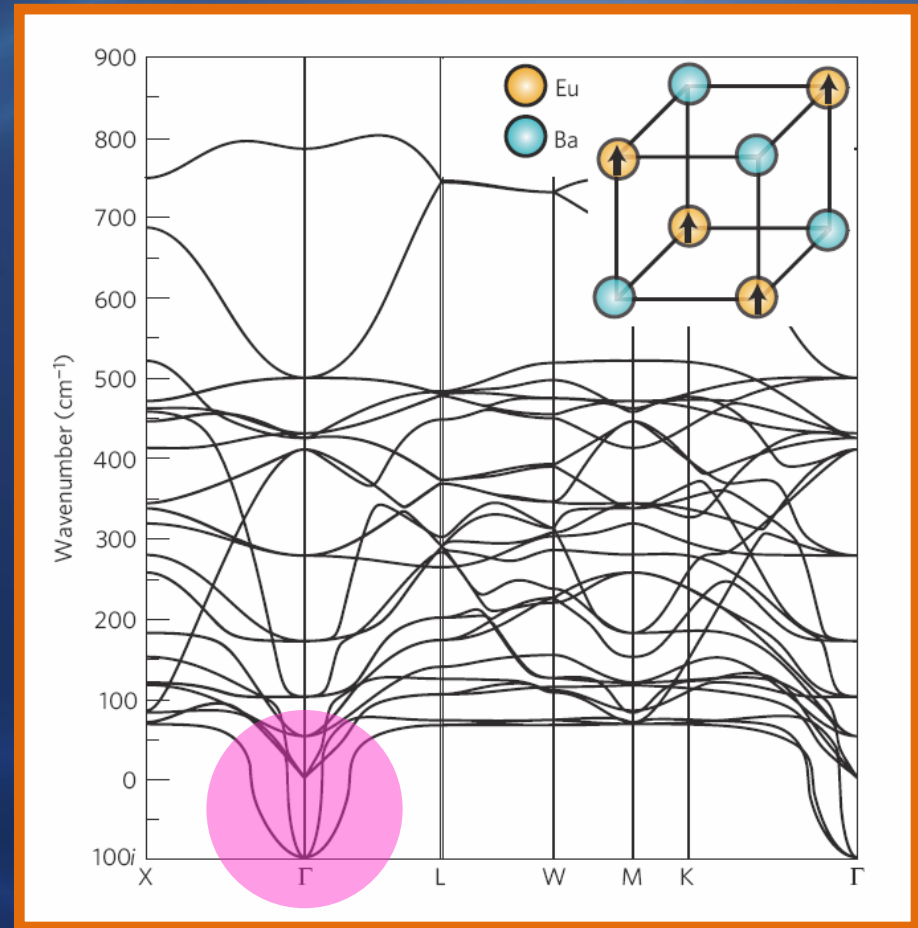
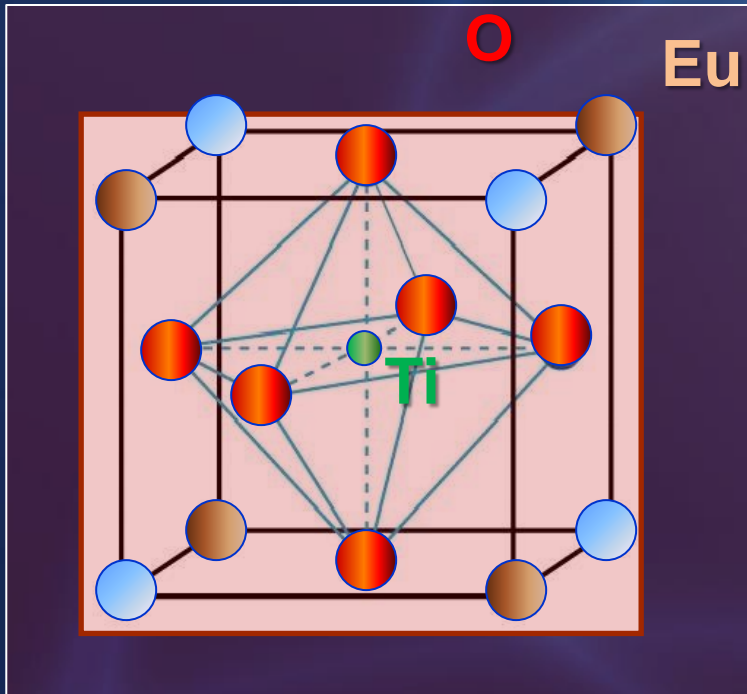
Ti

Ionic radius: **0.68 Å**

O

Ionic radius: **1.40 Å**

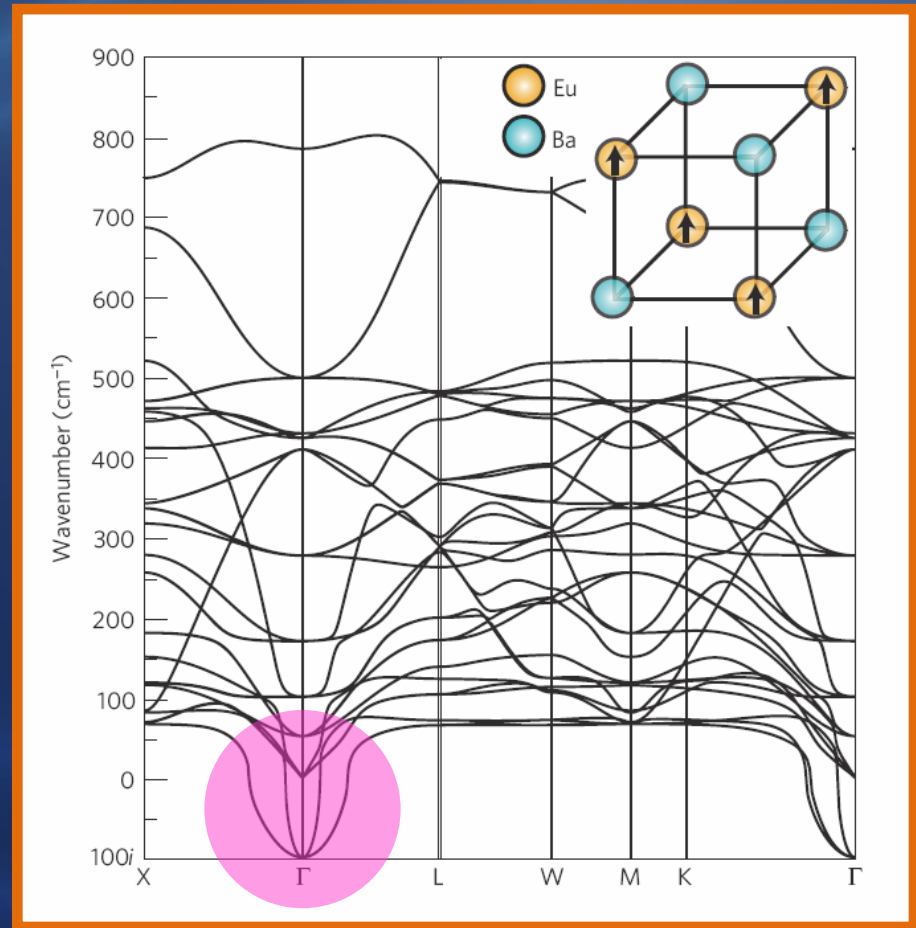
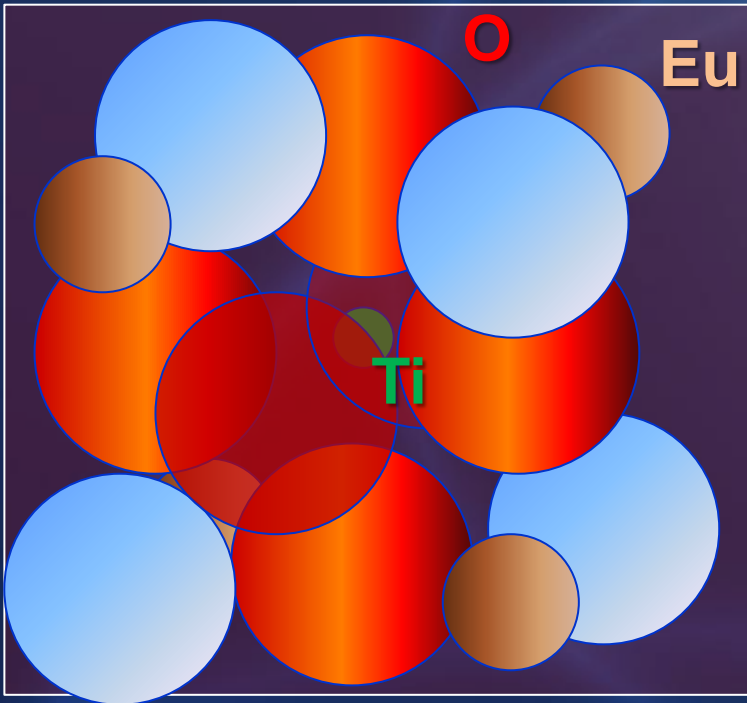
Calculated phonon dispersion of ferromagnetic $\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$ in its high-symmetry reference structure with pseudo-cubic lattice constant $a_0 = 3,95 \text{ \AA}$.



The imaginary-frequency polar phonon at Γ indicates a structural instability to a ferroelectric phase.

The inset shows the supercell of the ferromagnetic $\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$ ordered alloy used in our calculations. The Ti and O ions are omitted for clarity; arrows represent the Eu magnetic moments.

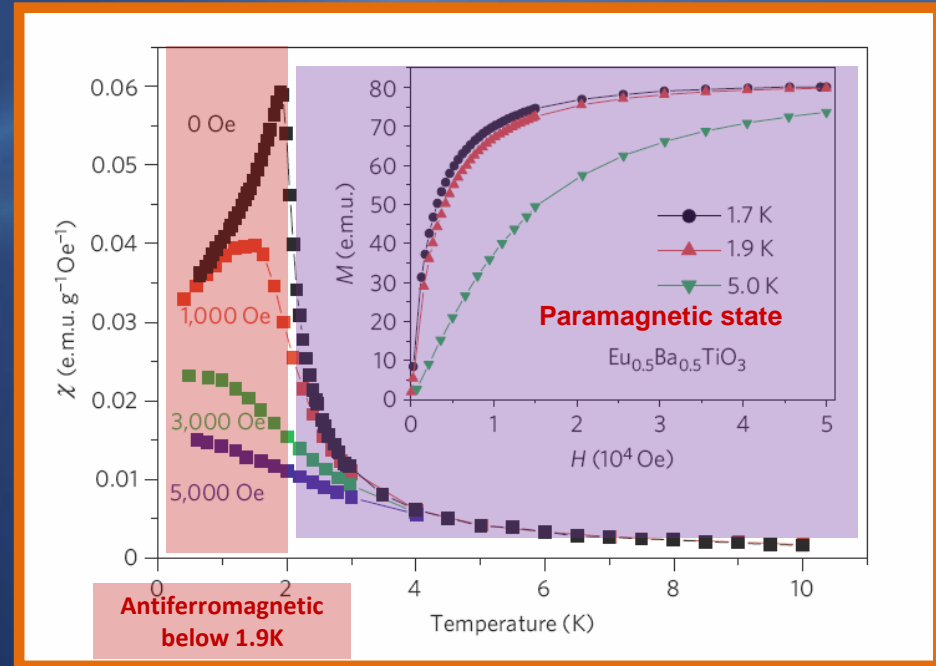
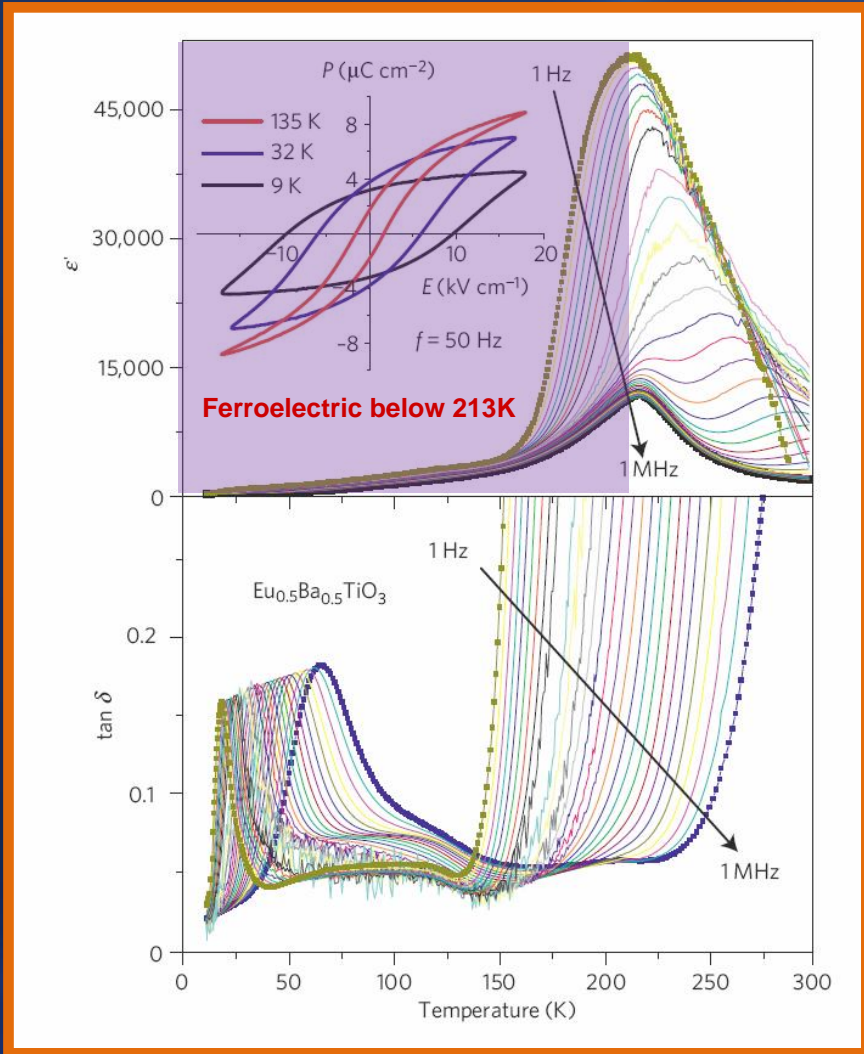
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Eu_{0.5}Ba_{0.5}TiO₃ ceramics



Temperature dependence of ac magnetic susceptibility, at various static magnetic fields and a frequency of 214 Hz.

The inset shows magnetization curves at various temperatures. We note that no hysteresis in magnetization was observed.

Temperature dependence of permittivity and dielectric loss in Eu_{0.5}Ba_{0.5}TiO₃ ceramics. The arrows indicate the direction of increasing frequency and the colours are for clarity to assist the eye in distinguishing the lines.

The inset shows ferroelectric hysteresis loops measured at three temperatures and 50 Hz.

In summary:

$\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$ ceramics confirmed their desirable ferroelectric polarization and absence of magnetic ordering above 1.9 K.

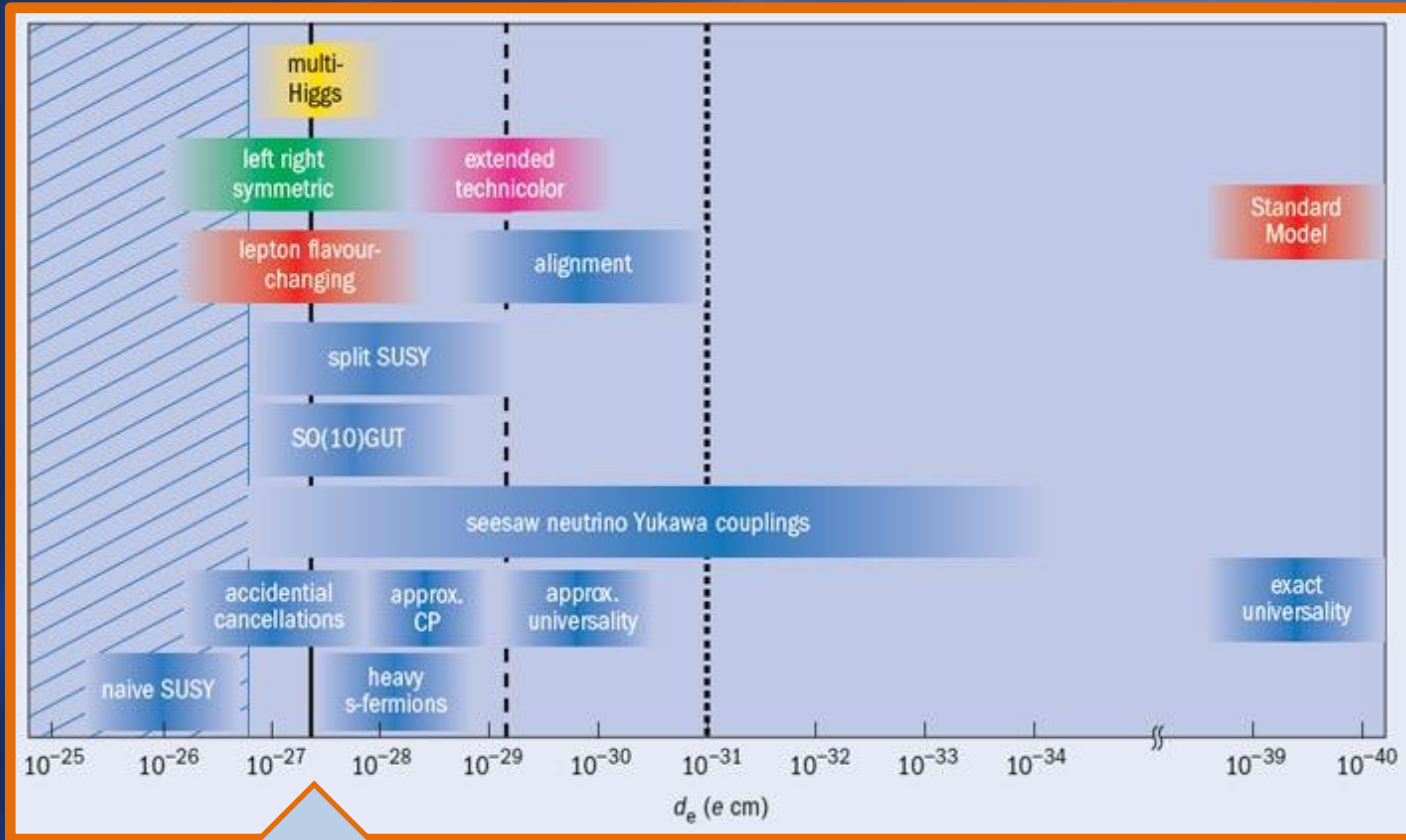
The search for the permanent dipole moment of the electron using $\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$ is now underway. Initial measurements have already achieved an EDM upper limit of $5 \cdot 10^{-23} \text{ e}\cdot\text{cm}$, which is within a factor of 10 of the *current record* with a solid-state-based EDM search.

A number of **systematic effects that may mask the EDM** signal is being controlled. The primary error originates from ferroelectric hysteresis-induced heating of the samples during polarization reversal. This **heating** gives rise to a change in magnetic susceptibility, which, in a non-zero external magnetic field, leads to an undesirable sample magnetization response. We are working to control the absolute magnetic field at the location of the samples to the **0.1 μG** level .

(Dla Ziemi indukcja magnetyczna $B = 30\mu\text{T} = 3 \cdot 10^{-3}\text{G} = 3\text{mG} = 3000 \mu\text{G}$)

Sensitivity of **$1 \cdot 10^{-28} \text{ e}\cdot\text{cm}$** should be achievable

Electron EDM predictions of various theories beyond the Standard Model



Experiments performed in 2002 at the University of California, Berkeley, showed that **if the electric dipole moment (EDM) of the electron exists, the value must be smaller than $10^{-27} e \text{ cm}$.**

Current maximum sensitivity

Next generation of cold-molecule experiments

Yale/Harvard experiment using thorium monoxide ThO



Czyż eEDM nie jest nieskończenie małą abstrakcją?



**Czyż eEDM nie jest
nieskończenie małą abstrakcją?**

